

A note about the torsion of null curves in the 3-dimensional Minkowski spacetime and the Schwarzian derivative

Zbigniew Olszak

Wrocław University of Technology, Institute of Mathematics and Computer Science, 50-370 Wrocław, Poland

Abstract. The main topic of this paper is to show that in the 3-dimensional Minkowski spacetime, the torsion of a null curve is equal to the Schwarzian derivative of a certain function appearing in a description of the curve. As applications, we obtain descriptions of the slant helices, and null curves for which the torsion is of the form $\tau = -2\lambda s$, s being the pseudo-arc parameter and $\lambda = \text{const} \neq 0$.

1. Introduction

There are very many papers about geometric properties of null curves in the Minkowski spacetimes. We refer the monographs [4, 5], and the survey articles [3, 11, 12], etc.

On the other hand, there is the classical notion of the Schwarzian derivative in mathematical analysis. This notion has many important applications in mathematical analysis (real and complex) and differential geometry; see [6, 7, 13–15], etc. The author is specially inspired by the paper [7], where it is shown a strict relation between the Schwarzian derivative and the curvature of worldlines in 2-dimensional Lorentzian manifolds of constant curvature.

In the presented short paper, we will show that the torsion of a null curve in the 3-dimensional Minkowski spacetime \mathbb{E}_1^3 is equal to the Schwarzian derivative of a certain function appearing in a description of the curve. Descriptions of the slant helices are obtained, and null curves for which the torsion is given by $\tau = -2\lambda s$, s being the pseudo-arc parameter and $\lambda = \text{const} \neq 0$.

2. Preliminaries

Let \mathbb{E}_1^3 be the 3-dimensional Minkowski spacetime, that is, the Cartesian \mathbb{R}^3 endowed with the standard Minkowski metric g given with respect to the Cartesian coordinates (x, y, z) by

$$g = dx \otimes dx + dy \otimes dy - dz \otimes dz, \quad (1)$$

or as the symmetric 2-form $g = dx^2 + dy^2 - dz^2$.

2010 *Mathematics Subject Classification.* Primary 53A35; Secondary 53B30, 53B50, 53C50.

Keywords. Minkowski spacetime of dimension 3; null curve; light-like curve; torsion of a curve; null helix; slant helix; Airy function; Schwarzian derivative.

Received: dd Month yyyy; Accepted: dd Month yyyy

Communicated by (name of the Editor, mandatory)

Email address: zbigniew.olszak@pwr.edu.pl (Zbigniew Olszak)

Let $\alpha: I \rightarrow \mathbb{E}_1^3$ be a null (ligh-like) curve in \mathbb{E}_1^3 , I being an open interval. Thus, $g(\alpha', \alpha') = 0$, that is, $g(\alpha'(t), \alpha'(t)) = 0$ for any $t \in I$. We also assume that the curve is non-degenerate, in the sense the three vector fields $\alpha', \alpha'', \alpha'''$ are linearly independent at every point of the curve.

Since $g(\alpha', \alpha') = 0$ and $g(\alpha', \alpha'') = 0$, it must be that $g(\alpha'', \alpha'') > 0$. A parametrization of the null curve is said to be pseudo-arc (or distinguished) if $g(\alpha'', \alpha'') = 1$. A null curve can always be parametrized by a pseudo-arc parameter. However, such a parameter is not uniquely defined. Precisely, for a null curve α , if s_1 is a pseudo-arc parameter, then s_2 is a pseudo-arc parameter if and only if there exists a constant c such that $s_2 = \pm s_1 + c$.

In the sequel, we assume that the parametrization of a null curve is pseudo-arc, and we denote such a parameter by s .

In the next section, we need the standard theorms concerning of null curves which can be formulated in the following manner (see e.g. [3, 5, 11, 12]):

Let α be a null curve in the 3-dimensional Minkowski spacetime \mathbb{E}_1^3 . Then, there exists the only one Cartan moving frame $(\mathbf{L} = \alpha', \mathbf{N}, \mathbf{W})$ and the function τ defined along the curve α and such that

$$g(\mathbf{L}, \mathbf{N}) = g(\mathbf{W}, \mathbf{W}) = 1, \quad g(\mathbf{L}, \mathbf{L}) = g(\mathbf{L}, \mathbf{W}) = g(\mathbf{N}, \mathbf{N}) = g(\mathbf{N}, \mathbf{W}) = 0, \quad (2)$$

and the following system of differential equations

$$\mathbf{L}' = \mathbf{W}, \quad \mathbf{N}' = \tau \mathbf{W}, \quad \mathbf{W}' = -\tau \mathbf{L} - \mathbf{N} \quad (3)$$

is satisfied. These vector fileds are given by

$$\mathbf{L} = \alpha', \quad \mathbf{W} = \alpha'', \quad \mathbf{N} = -\alpha''' - \frac{1}{2}g(\alpha''', \alpha''')\alpha', \quad (4)$$

and the function τ by

$$\tau = \frac{1}{2}g(\alpha''', \alpha'''). \quad (5)$$

From these results it can be deduced that a given function τ on an open interval I , there exists the only one null curve $\alpha: I \rightarrow \mathbb{E}_1^3$ realizing (2) and (3) up to the orientation of this curve and up to the isometries of the Minkowski space \mathbb{E}_1^3 .

The triple $(\mathbf{L}, \mathbf{N}, \mathbf{W})$ defined in (4) is called the Frenet frame, the function τ defined in (5) is called the torsion, and the equations (3) are called the Frenet equations of the null curve α . Since

$$\det[\mathbf{L}, \mathbf{N}, \mathbf{W}] = \det[\alpha', \alpha'', \alpha'''], \quad (6)$$

the frames $(\mathbf{L}, \mathbf{N}, \mathbf{W})$ and $(\alpha', \alpha'', \alpha''')$ have the same orientations.

In the following section, we are going to expresse the torsion τ and the frame $(\mathbf{L}, \mathbf{N}, \mathbf{W})$ with the help of a special function related to a pseudo-arc parametrization of a null curve in \mathbb{E}_1^3 .

3. A description of the torsion

Let $\alpha: I \rightarrow \mathbb{E}_1^3$ be a null curve. Simplifying denotations, we write $\alpha(s) = (x(s), y(s), z(s))$, $s \in I$, where s is a pseudo-arc parameter, and $x(s)$, $y(s)$, $z(s)$ are certain functions of s . Then, we have

$$\alpha' = x' \frac{\partial}{\partial x} \Big|_{\alpha} + y' \frac{\partial}{\partial y} \Big|_{\alpha} + z' \frac{\partial}{\partial z} \Big|_{\alpha}.$$

For simplicity, instead of that, we will write $\alpha' = (x', y', z')$. And in the similar manner, the next derivatives of α will be written, e.g., $\alpha'' = (x'', y'', z'')$.

Using (1), our two assumptions: $g(\alpha', \alpha') = 0$ (the nullity condition), and $g(\alpha'', \alpha'') = 1$ (the pseudo-arc parametrization) give the following two equalities

$$x'^2 + y'^2 - z'^2 = 0, \quad (7)$$

$$x''^2 + y''^2 - z''^2 = 1. \quad (8)$$

One notes that the shapes of the equalities (7) and (8) exclude the situation when at least one of the functions x' , y' , z' vanishes on an open subinterval of I . In the sequel, restricting slightly the assumptions, we will consider only the case when $x' \neq 0$, $y' \neq 0$ and $z' \neq 0$ on I .

It is a standard and elementary idea that from (7), it follows that

$$x' = h, \quad y' = \frac{h}{2} \left(f - \frac{1}{f} \right), \quad z' = \frac{h}{2} \left(f + \frac{1}{f} \right), \quad (9)$$

f and h being certain non-zero functions on I . Hence,

$$\begin{aligned} x'' &= h', \\ y'' &= \frac{fh'(f^2 - 1) + hf'(f^2 + 1)}{2f^2}, \\ z'' &= \frac{fh'(f^2 + 1) + hf'(f^2 - 1)}{2f^2}. \end{aligned}$$

In view of the above relations, the equality (8) turns into $h^2 f'^2 = f^2$. Hence, f' is non-zero (and has constant sign) on I . Consequently,

$$h = \varepsilon \frac{f}{f'}, \quad \varepsilon = \pm 1.$$

Thus, for the vector field \mathbf{L} (cf. (4)), we have

$$\mathbf{L} = \alpha' = \frac{\varepsilon}{2f'} (2f, f^2 - 1, f^2 + 1). \quad (10)$$

Consequently, we get the following description of the curve α

$$\alpha(s) = \alpha(s_0) + \frac{\varepsilon}{2} \int_{s_0}^s \frac{1}{f'(t)} (2f(t), f^2(t) - 1, f^2(t) + 1) dt, \quad s, s_0 \in I.$$

Conversely, if a curve α is given by the last formula, then (7) and (8) are fulfilled so that the curve is null and not geodesic, and the parameter s is distinguish.

From (10), we obtain for the vector field \mathbf{W} (cf. (4)),

$$\mathbf{W} = \alpha'' = -\frac{\varepsilon f''}{2f'^2} (2f, f^2 - 1, f^2 + 1) + \varepsilon (1, f, f). \quad (11)$$

From (11), we find

$$\alpha''' = \varepsilon \frac{2f''^2 - f'f'''}{2f'^3} (2f, f^2 - 1, f^2 + 1) - \frac{\varepsilon f''}{f'} (1, f, f) + \varepsilon f' (0, 1, 1). \quad (12)$$

To compute $g(\alpha''', \alpha''')$, using (1), we find at first the following

$$\begin{aligned} g((2f, f^2 - 1, f^2 + 1), (2f, f^2 - 1, f^2 + 1)) &= 0, \\ g((2f, f^2 - 1, f^2 + 1), (1, f, f)) &= 0, \quad g((2f, f^2 - 1, f^2 + 1), (0, 1, 1)) = -2, \\ g((1, f, f), (1, f, f)) &= 1, \quad g((1, f, f), (0, 1, 1)) = 0, \quad g((0, 1, 1), (0, 1, 1)) = 0. \end{aligned}$$

Then, having (12) and applying the above formulas, we get

$$g(\alpha''', \alpha''') = \frac{2f'f''' - 3f''^2}{f'^2}. \quad (13)$$

In view of (13) and (5), the torsion must be of the form

$$\tau = \frac{2f'f''' - 3f''^2}{2f'^2} = \left(\frac{f''}{f'}\right)' - \frac{1}{2}\left(\frac{f''}{f'}\right)^2. \quad (14)$$

Now, it is important to note that the right hand side of the formula (14) is just the Schwarzian derivative of the function f , which is usually denoted by $S(f)$. Thus, $\tau = S(f)$.

Finally, applying (10), (12) and (13) into (4), we find the vector field

$$\mathbf{N} = -\frac{\varepsilon f''^2}{4f'^3} (2f, f^2 - 1, f^2 + 1) + \frac{\varepsilon f''}{f'} (1, f, f) - \varepsilon f'(0, 1, 1). \quad (15)$$

Summarizing the above considerations, we can formulate the following theorem.

Theorem 1. *Let \mathbb{E}_1^3 be the 3-dimensional Minkowski spacetime. Any (non-degenerate) null curve α in \mathbb{E}_1^3 can be parametrized in the following way*

$$\alpha(s) = \alpha(s_0) + \frac{\varepsilon}{2} \int_{s_0}^s \frac{1}{f'(t)} (2f(t), f^2(t) - 1, f^2(t) + 1) dt, \quad s, s_0 \in I, \quad (16)$$

where s is a pseudo-arc parameter, I is a certain open interval, f is a non-zero function with non-zero derivative f' on I . The torsion τ of such a curve is equal to the Schwarzian derivative of the function f , that is,

$$\tau = S(f) = \left(\frac{f''}{f'}\right)' - \frac{1}{2}\left(\frac{f''}{f'}\right)^2. \quad (17)$$

The vector fields forming the Frenet frame of the curve α are given by the formulas (10), (11) and (15).

Remark 1. Applying formulas (10), (11), (15), it can be verified that

$$\det[\mathbf{L}, \mathbf{N}, \mathbf{W}] = \varepsilon.$$

This together with (6) implies that the constant ε appearing in (16) corresponds to the orientation of the curve α . Note that the torsion does not depend on the orientation of the curve. Moreover, the torsion and the orientation does not depend on the sign of the function f .

Remark 2. The Schwarzian derivative S is an invariant of a fractional-linear transformation T of the 1-dimensional real projective space $\mathbb{RP}^1 = \mathbb{R} \cup \infty$ (cf. e.g. [13]). That is, $S(T \circ f) = S(f)$ if f is a function on \mathbb{RP}^1 and $T: \mathbb{RP}^1 \rightarrow \mathbb{RP}^1$ is given by

$$T(r) = \frac{ar + b}{cr + d}, \quad r \in \mathbb{RP}^1, \quad a, b, c, d \in \mathbb{R}, ad - bc \neq 0. \quad (18)$$

We can apply the above fact seeking for null curves with given torsion τ . However, we should be careful since the domains of our functions f and $T \circ f$ may be defined only on some open subintervals lying on the real line \mathbb{R} .

4. Null Cartan helices

It is well-known that there are exactly three types of null curves with constant torsion in the Minkowski spacetime \mathbb{E}_1^3 (cf e.g., [9]) up to the orientation of the curve and up to the isometries of the space. They are often called the null Cartan helices.

As a first application of the results from the previous section, we demonstrate how these classes of curves can be recovered from their torsions.

(a) For $f(s) = s$, it holds $S(f) = 0$. In (16), we put $f(s) = s$, $s_0 = 0$, $\alpha(s_0) = (0, 0, 0)$, $\varepsilon = 1$. Then, we obtain the curve

$$\alpha(s) = \frac{1}{6} (3s^2, s^3 - 3s, s^3 + 3s),$$

for which by (17) we have $\tau = 0$. Thus, α is a positively oriented null Cartan helix of zero torsion.

(b) For $f(s) = -\cot(cs/2)$, it holds $S(f) = c^2/2$. In (16), we put

$$f(s) = -\cot \frac{cs}{2}, \quad \alpha(0) = \left(\frac{1}{c^2}, 0, 0 \right), \quad \varepsilon = 1, \quad c = \text{const.} > 0.$$

Then, we obtain the curve

$$\alpha(s) = \frac{1}{c^2} (\cos(cs), \sin(cs), cs),$$

for which by (17) it holds $\tau = c^2/2$. Thus, α is a positively oriented null Cartan helix of constant positive torsion.

(c) For $f(s) = e^{cs}$, it holds $S(f) = -c^2/2$. In (16), we put

$$f(s) = e^{cs}, \quad \alpha(0) = \left(0, \frac{1}{c^2}, 0 \right), \quad \varepsilon = 1, \quad c = \text{const.} > 0.$$

Then, we obtain the curve

$$\alpha(s) = \frac{1}{c^2} (cs, \cosh(cs), \sinh(cs)),$$

for which by (17) we have $\tau = -c^2/2$. Thus, α is a positively oriented null Cartan helix of constant negative torsion.

Thus, we have seen the following:

Corollary 1. *Null helices in \mathbb{E}_1^3 form the three classes described in (a) – (c) in the above. The description is valid up to the pseudo-arc parameter changes, up to the orientation of the curve, and up to the isometries of the space.*

A curve $\alpha: I \rightarrow \mathbb{E}_1^3$ is called a general (or generalized) helix if there exists a non-zero vector V in \mathbb{E}_1^3 such that $g(\alpha', V) = \text{const.}$; cf. [8, 9, 17], etc. This means that tangent indicatrix is laid in a plane or, equivalently, there exists a non-zero constant vector V in \mathbb{E}_1^3 for which $g(\alpha'', V) = 0$, that is, V is orthogonal to the acceleration vector field α'' .

For null curves, it is already proved that null general helices in \mathbb{E}_1^3 are precisely the null Cartan helices; cf. ibidem.

5. Null slant helices

Following the ideas of [1, 2, 10], a slant helix is defined to be the curve (null as well as non-null) in \mathbb{E}_1^3 which satisfies the condition

$$g(\alpha'', V) = c = \text{const.} \tag{19}$$

along the curve α , where V is a constant vector. Thus, a general helix is a slant helix with $c = 0$. Conversely, a slant helix with $c = 0$ becomes a general helix. In [1, Theorem 1.4], it is proved that a null curve in \mathbb{E}_1^3 is a slant helix if and only if its torsion is given by

$$\tau = \frac{a}{(cs + b)^2}, \quad a, b, c = \text{const.}, \quad (20)$$

where c is just the constant realizing (19).

As the second applications of the results from Section 3, we will describe the null slant helices in \mathbb{E}_1^3 which are different from the usual helices ($a \neq 0$ and $c \neq 0$ in (20)).

Note that moving the pseudo-arc parameter s into $s - b/c$ and next modifying slightly the constant a , we can write the condition (20) as

$$\tau = \frac{a}{2s^2}, \quad a = \text{const} \neq 0. \quad (21)$$

We can also assume that $s > 0$. Using (2) and (3), it can be checked that when the relation (21) is fulfilled, then for the vector

$$V = -\frac{a}{2s}\mathbf{L} + s\mathbf{N} + \mathbf{W}$$

it holds $V' = 0$ and $g(\alpha'', V) = g(\mathbf{W}, V) = 1$ (cf. ibidem).

(a) In (16), we put

$$f(s) = \ln s, \quad s_0 = 1, \quad \alpha(s_0) = \frac{1}{8}(-2, -1, 3), \quad \varepsilon = 1.$$

Then, we obtain the curve

$$\alpha(s) = \frac{s^2}{8}(2(2 \ln s - 1), 2 \ln^2 s - 2 \ln s - 1, 2 \ln^2 s - 2 \ln s + 3),$$

for which by (17) it holds

$$\tau = S(f) = \frac{1}{2s^2}.$$

Thus, α is a slant helix realizing (21) with $a = 1$.

(b) Let $a > 1$ and $b = \sqrt{a-1} > 0$. In (16), we put

$$f(s) = \tan\left(\frac{1}{2} \ln s^b\right), \quad s_0 = 1, \quad \alpha(s_0) = \frac{1}{b}\left(-\frac{b}{b^2+4}, -\frac{2}{b^2+4}, \frac{1}{2}\right), \quad \varepsilon = 1.$$

Then, we obtain the curve

$$\alpha(s) = \frac{s^2}{b}\left(\frac{2 \sin(\ln s^b) - b \cos(\ln s^b)}{b^2+4}, -\frac{2 \cos(\ln s^b) + b \sin(\ln s^b)}{b^2+4}, \frac{1}{2}\right),$$

for which by (17) it holds

$$\tau = S(f) = \frac{1+b^2}{2s^2} = \frac{a}{2s^2}.$$

Thus, α is a slant helix realizing (21) with $a > 1$.

(c) Let $0 \neq a < 1$. Then for $b = \sqrt{1-a}$, we have $b > 0$ and $b \neq 1$. Consider the case $a \neq -3$, that is, $b \neq 2$. In (16), we put

$$f(s) = s^{-b}, \quad s_0 = 1, \quad \alpha(s_0) = \frac{1}{2b}\left(-1, \frac{2b}{b^2-4}, \frac{4}{b^2-4}\right), \quad \varepsilon = 1.$$

Then, we obtain the curve

$$\alpha(s) = \frac{s^2}{2b} \left(-1, \frac{s^{-b}}{b-2} + \frac{s^b}{b+2}, \frac{s^{-b}}{b-2} - \frac{s^b}{b+2} \right),$$

for which by (17) it holds

$$\tau = S(f) = \frac{1-b^2}{2s^2} = \frac{a}{2s^2}.$$

Thus, α is a slant helix realizing (21) with $-3 \neq a < 1$.

(d) In (16), we put

$$f(s) = \frac{1}{s^2}, s_0 = 1, \alpha(s_0) = \frac{1}{16}(-4, 1, -1), \varepsilon = 1.$$

Then, we obtain the curve

$$\alpha(s) = \frac{1}{16}(-4s^2, s^4 - 4 \ln s, -s^4 - 4 \ln s),$$

for which by (17) it holds

$$\tau = S(f) = -\frac{3}{2s^2}.$$

Thus, α is a slant helix realizing (21) with $a = -3$.

Thus, we have shown the following:

Corollary 2. *Null slant helices in \mathbb{E}_1^3 form the four classes described in (a) – (d) in the above. The description is valid up to the pseudo-arc parameter changes, up to the orientation of the curve, and up to the isometries of the space.*

6. Null curves with the torsion proportional to the pseudo-arc parameter

In this section, we determine the null curves in \mathbb{E}_1^3 for which $\tau = -2\lambda s$, $\lambda = \text{const.} \neq 0$. We will use the formula (17).

According to our Theorem, we need at first to find a solution of the differential equation

$$\left(\frac{f''}{f'} \right)' - \frac{1}{2} \left(\frac{f''}{f'} \right)^2 = -2\lambda s. \quad (22)$$

We seek for solutions of this equation in the form

$$f(s) = \int \frac{ds}{\phi^2(s)}, \quad (23)$$

ϕ being an unknown function. Then the equation (22) becomes the following differential equation

$$\phi'' - \lambda s \phi = 0. \quad (24)$$

The general solution of the above equation is

$$\phi(s) = c_1 \text{Ai}(\mu s) + c_2 \text{Bi}(\mu s), \mu = \sqrt[3]{\lambda}, c_1, c_2 = \text{const.}$$

where Ai and Bi are the Airy functions of the first and second kind, respectively. For the solutions of (24) and for the special Airy functions, we refer [16], [18], etc. In the below calculations, we use the basic properties of these functions.

For our purpose, we take the only one solution of (24), say $\phi(s) = \text{Ai}(\mu s)$. Then, from (23) we get

$$f(s) = \frac{\pi}{\mu} \cdot \frac{\text{Bi}(\mu s)}{\text{Ai}(\mu s)}. \quad (25)$$

Next,

$$f'(s) = \frac{1}{\text{Ai}^2(\mu s)}. \quad (26)$$

Having (10) with $\varepsilon = 1$, and using (25) and (26), we can write α' as

$$\alpha'(s) = \left(\frac{\pi}{\mu} \text{Ai}(\mu s) \text{Bi}(\mu s), \frac{1}{2\mu^2} (\pi^2 \text{Bi}^2(\mu s) - \mu^2 \text{Ai}^2(\mu s)), \frac{1}{2\mu^2} (\pi^2 \text{Bi}^2(\mu s) + \mu^2 \text{Ai}^2(\mu s)) \right).$$

The integration of the last equality gives the following curve

$$\begin{aligned} \alpha(s) = & \left(\frac{\pi}{\mu^2} (\mu s \text{Ai}(\mu s) \text{Bi}(\mu s) - \text{Ai}'(\mu s) \text{Bi}'(\mu s)), \right. \\ & \frac{1}{2\mu^3} (\pi^2 (\mu s \text{Bi}^2(\mu s) - \text{Bi}'^2(\mu s)) - \mu^3 s \text{Ai}^2(\mu s) + \mu^2 \text{Ai}'^2(\mu s)), \\ & \left. \frac{1}{2\mu^3} (\pi^2 (\mu s \text{Bi}^2(\mu s) - \text{Bi}'^2(\mu s)) + \mu^3 s \text{Ai}^2(\mu s) - \mu^2 \text{Ai}'^2(\mu s)) \right), \end{aligned} \quad (27)$$

if the the initial condition at $s_0 = 0$ is

$$\alpha(0) = \frac{1}{2\sqrt[3]{9}\mu^3\Gamma(\frac{1}{3})} (2\sqrt{3}\mu\pi, \mu^2 - 3\pi^2, -\mu^2 - 3\pi^2).$$

Thus, we can formulate the following:

Corollary 3. *Null curves in \mathbb{E}_1^3 for which $\tau = -2\lambda s$, $\lambda = \text{const.} \neq 0$, are given by the formula (27) with $\mu = \sqrt[3]{\lambda}$ up to the pseudo-arc parameter changes, up to the orientation of the curve, and up to the isometries of the space.*

References

- [1] A. T. Ali and R. López, Slant helices in Minkowski space \mathbb{E}_1^3 , Journal of the Korean Mathematical Society 48 (2011) 159–167.
- [2] J. H. Choi and Y. H. Kim, Note on null helices in \mathbb{E}_1^3 , Bulletin of the Korean Mathematical Society 50 (2013) 885–899.
- [3] K. L. Duggal, A report on canonical null curves and screen distributions for lightlike geometry, Acta Applicandae Mathematicae 95 (2007) 135–149.
- [4] K. L. Duggal and A. Bejancu, Lightlike submanifolds of semi-Riemannian manifolds and applications, Kluwer Academic Publishers, Dordrecht, 1996.
- [5] K. L. Duggal and D. H. Jin, Null curves and hypersurfaces of semi-Riemannian manifolds, World Scientific Publishing, Hackensack, 2007.
- [6] C. Duval and L. Guieu, The Virasoro group and Lorentzian surfaces: the hyperboloid of one sheet, Journal of Geometry and Physics 33 (2000) 103–127.
- [7] C. Duval and V. Ovsienko, Lorentzian worldlines and the Schwarzian derivative, Functional Analysis and its Applications 34 (2000), No. 2, 135–137; translation from Funktsionalnyj Analiz i Ego Prilozheniya 34 (2000), No. 2, 69–72.
- [8] A. Ferrández, A. Giménez and P. Lucas, Null helices in Lorentzian space forms, International Journal of Modern Physics A 16 (2001) 4845–4863.
- [9] A. Ferrández, A. Giménez and P. Lucas, Null generalized helices in Lorentz-Minkowski spaces, Journal of Physics A: Mathematical and General 35 (2002) 8243–8251.
- [10] F. Gökçelik and İ. Gök, Null W -slant helices in \mathbb{E}_1^3 , Journal of Mathematical Analysis and Applications 420 (2014) 222–241.
- [11] J.-I. Inoguchi and S. Lee, Null curves in Minkowski 3-space, International Electronic Journal of Geometry 1 (2008), No. 2, 40–83.
- [12] R. Lopez, Differential geometry of curves and surfaces in Lorentz-Minkowski space, International Electronic Journal of Geometry 7 (2014), No. 1, 44–107.

- [13] B. Osgood, Old and new on the Schwarzian derivative, In: Duren, Peter (ed.) et al., Quasiconformal mappings and analysis. Proceedings of the international symposium, Ann Arbor, MI, USA, August 1995, pp. 275-308, New York, NY: Springer, 1998.
- [14] B. Osgood and D. Stowe, *The Schwarzian derivative and conformal mapping of Riemannian manifolds*, Duke Mathematical Journal 67 (1992) 57–99.
- [15] V. Ovsienko and S. Tabachnikov, *What is the Schwarzian derivative?*, Notices of the American Mathematical Society 56 (2009) 34–36.
- [16] A. D. Polyanin and V. F. Zaitsev, *Handbook of exact solutions for ordinary differential equations*, 2nd Edition, Chapman & Hall/CRC, Boca Raton, 2003.
- [17] B. Şahin, E. Kiliç and R. Güneş, *Null helices in \mathbb{R}^3* , Differential Geometry - Dynamical Systems 3 (2001), No. 2, 31–36.
- [18] O. Vallée and M. Soares, *Airy functions and applications to physics*, Imperial College Press, London, 2004.